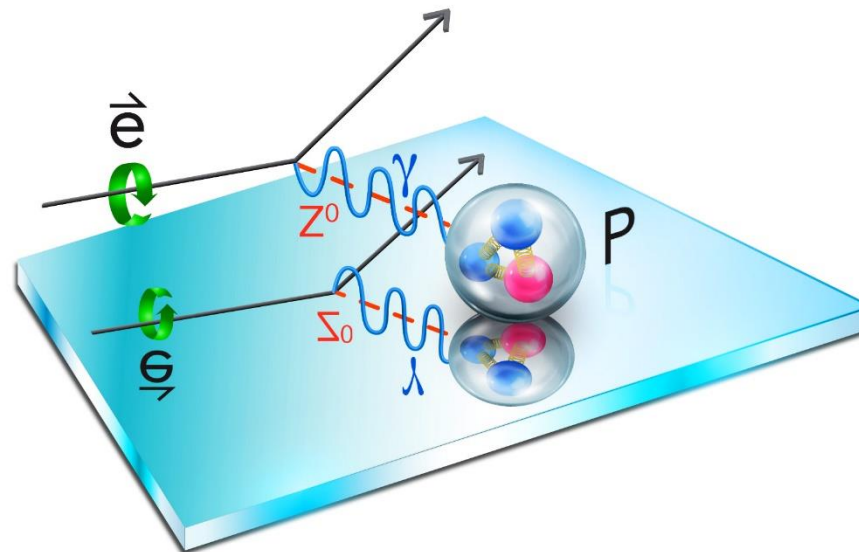
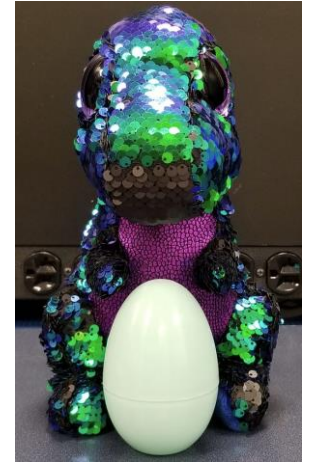
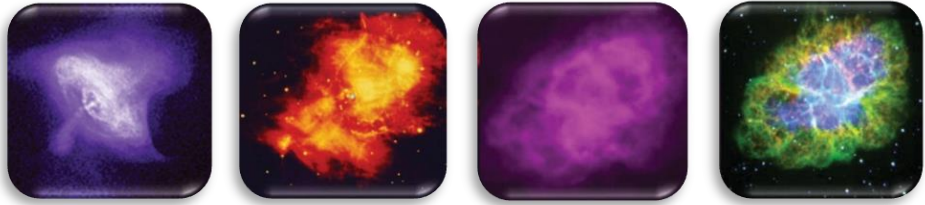


PREX/CREX Overview

Juliette Mammei
for the PREX and CREX Collaborations



Connecting heaven and earth

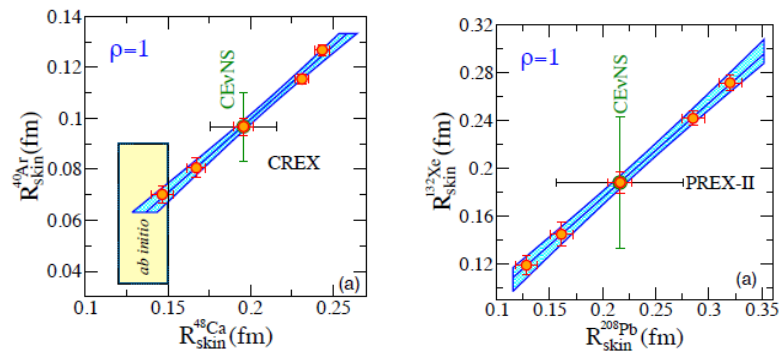


Crab Nebula (X-ray, infrared, radio, visible)

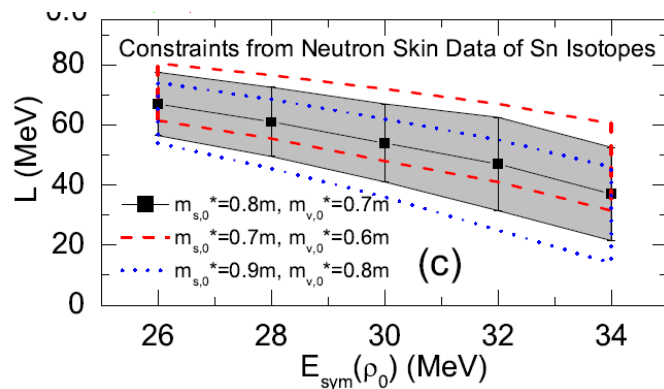


Gravitational (LIGO/Virgo)

It all has to hang together...



Piekarewicz (CEvNS2019)
coherent neutrino scattering



June 2021

Chen et al.

<https://arxiv.org/abs/1004.4672v2>

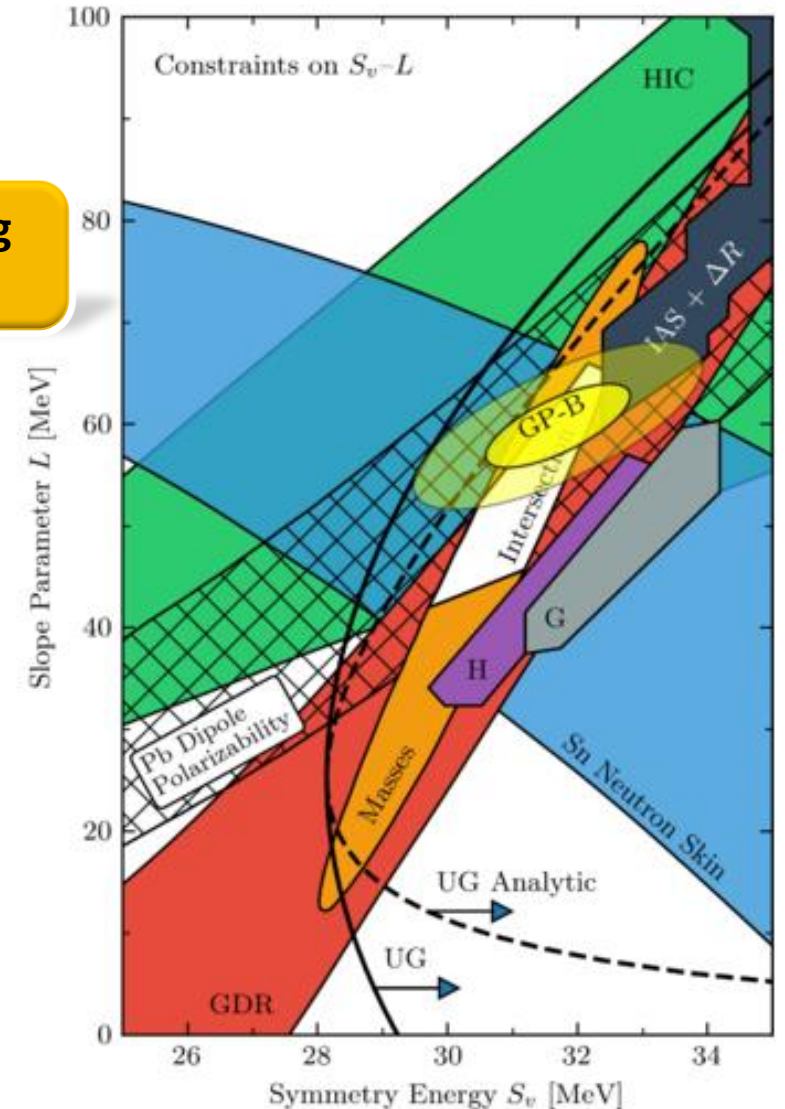
Earth-based (nuclei)

Heavy ion collisions
GDR
neutron skin in Tin isotopes
nuclear masses
Dipole polarizability
Coherent neutrino scattering

Space-based (neutron stars)

NICER – EM
LIGO/Virgo – tidal deformability

Same particles, same EOS,
18 orders of magnitude different in size!



Drischler et al.

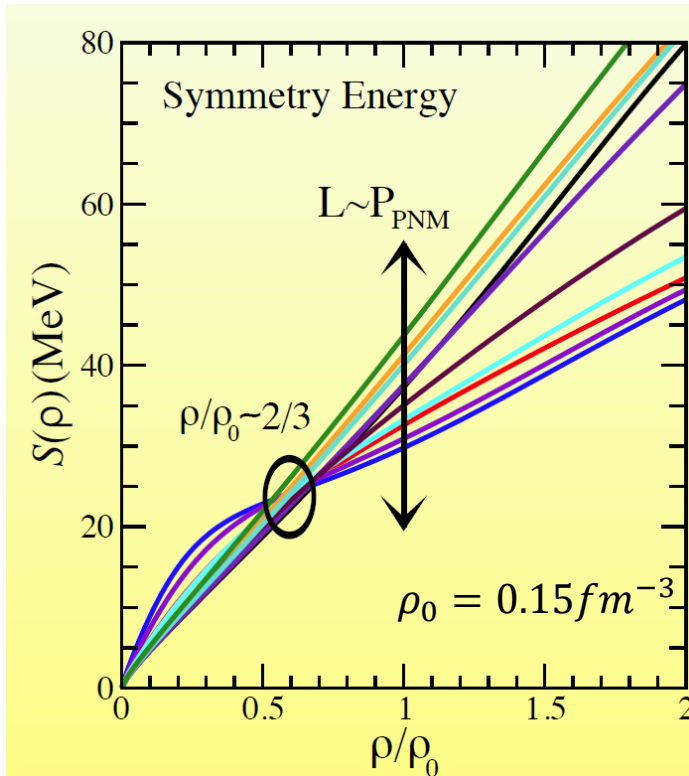
<https://arxiv.org/abs/2004.07232>

Equation of State of Neutron Matter

E/A for symmetric nuclear matter Symmetry energy

$$\underbrace{\varepsilon(\rho, \alpha = 1)}_{\text{EOS of pure neutron matter}} \approx \underbrace{\varepsilon_0(\rho)}_{\text{E/A for symmetric nuclear matter}} + \underbrace{\alpha^2 S(\rho)}_{\text{Symmetry energy}} \approx \epsilon_0 + Lx$$

EOS of pure neutron matter



$$\alpha = (N - Z)/A$$

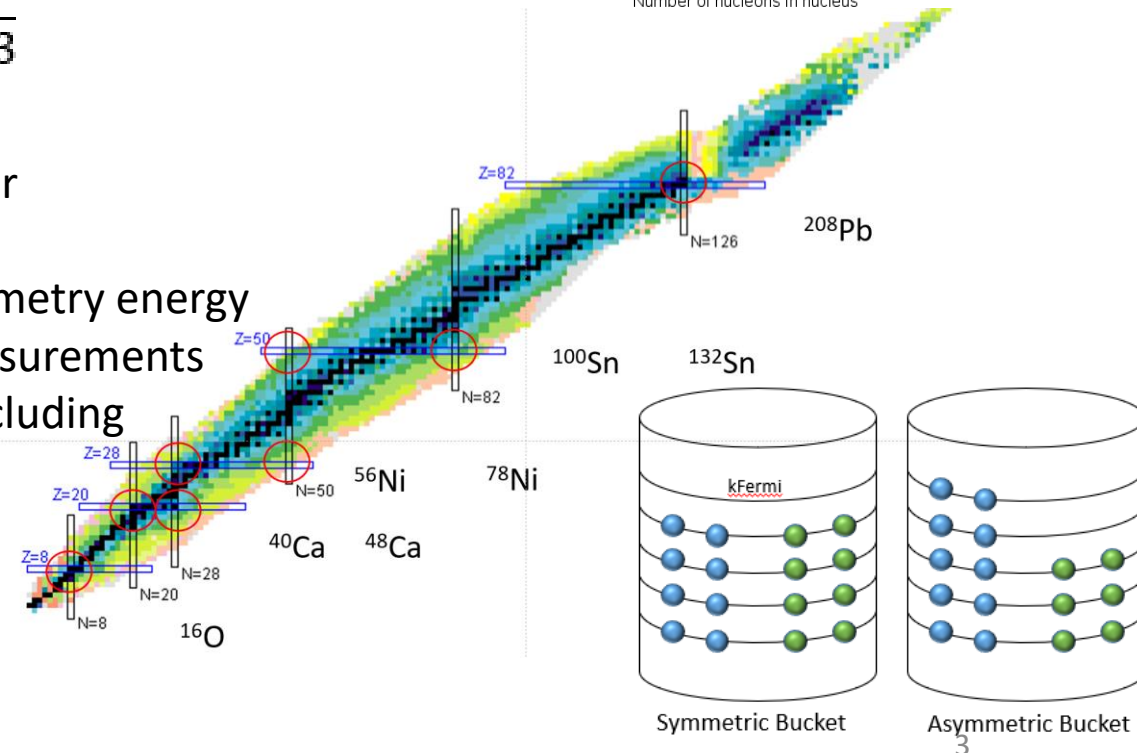
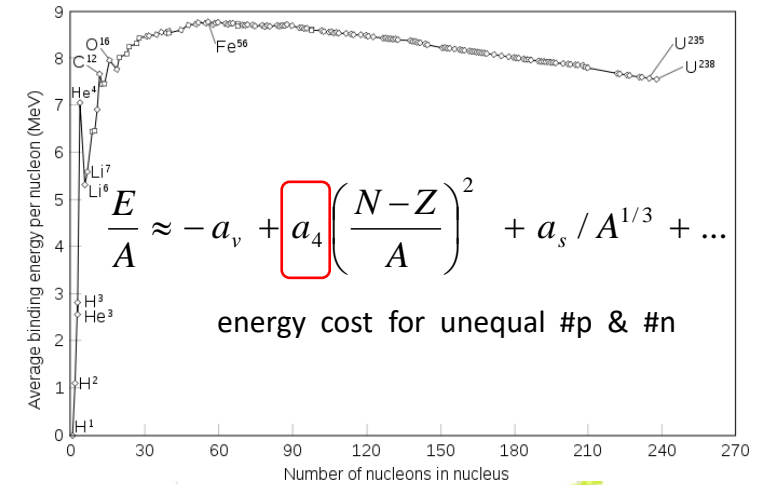
$$x = (\rho - \rho_0)/3\rho_0$$

$$P_0 = \left(\rho^2 \frac{\partial \varepsilon_{PNM}}{\partial \rho} \right) \Big|_{\rho = \rho_0} \approx \rho_0 \frac{L}{3}$$

Pressure of pure neutron matter

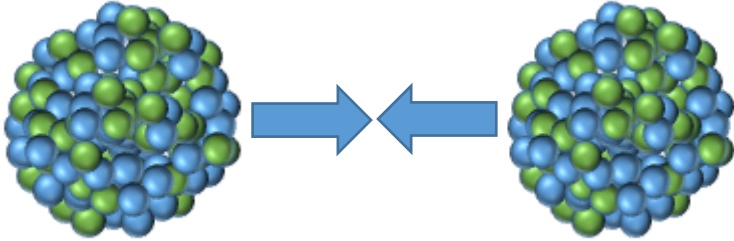
L – density dependence of symmetry energy unconstrained by isoscalar measurements (~symmetric nuclear matter, including nuclear binding energies, etc.)

Doubly-magic and neutron rich nuclei to simplify corrections

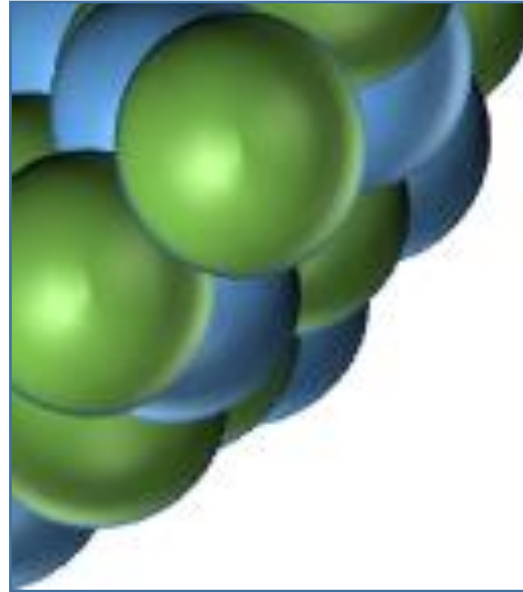


Parity-violating electron scattering (PVES)

other methods, like HIC, have strong interaction uncertainties

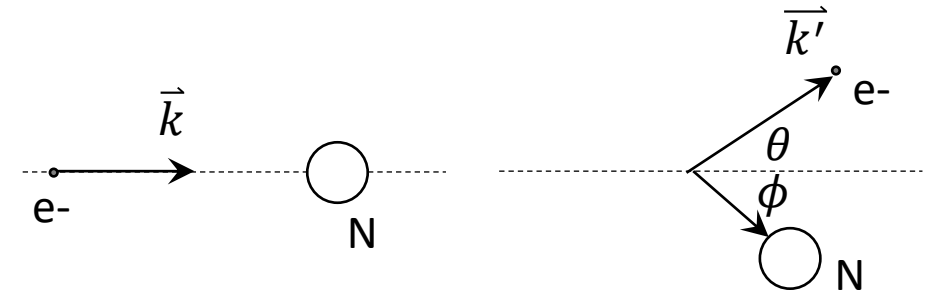
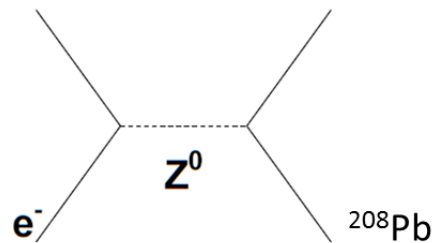
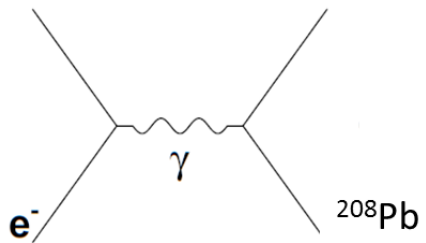


in electron scattering the probe doesn't interact via the strong force



Electrons with different helicities "see" different potentials for the nucleus because of parity-violation in the weak interaction

does interact via the E&M and weak forces



$$q^2 = (\vec{k}' - \vec{k})^2$$

Elastic scattering
 $\Rightarrow -q^2 = Q^2 = 4EE' \sin^2 \theta$

Neutron skin with PVES

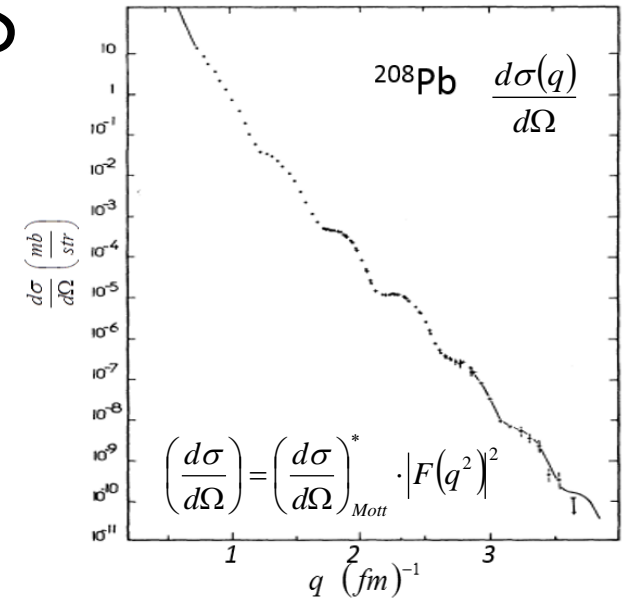
$$A_{PV} = \frac{\sigma_+ - \sigma_-}{\sigma_+ + \sigma_-} \approx \frac{\text{[Feynman diagrams: } e^- \gamma 208\text{Pb} + e^- Z^0 208\text{Pb} \text{]} }{2 \cdot \text{[Feynman diagram: } e^- \gamma 208\text{Pb} \text{]}}$$

$$\approx \frac{G_F Q^2 |Q_W| F_W(Q^2)}{4\sqrt{2}\pi\alpha Z F_{ch}(Q^2)}$$

The Fourier transform of the weak “form factor” $F_W(Q^2)$ gives the weak charge density as a function of radius, just as it does for the charge form factor

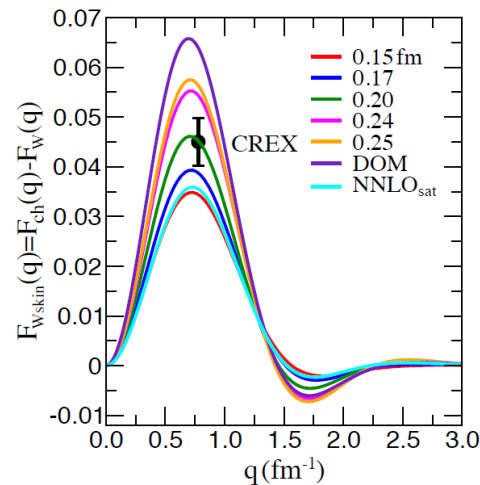
$$Q_{weak}^p = 1 - 4\sin^2 \theta_W \approx 0$$

$$Q_{weak}^n = -1$$

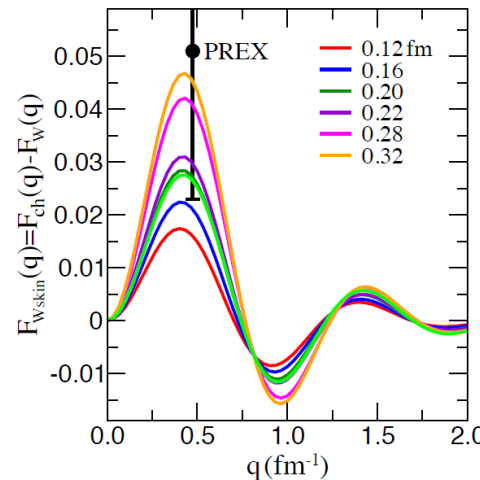


Measurement of $F_n(Q^2)$ at a single Q^2 translates to a measurement of R_n via mean-field nuclear models

At low Q^2 there is a tight correlation between R_n and $F_n(Q^2)$

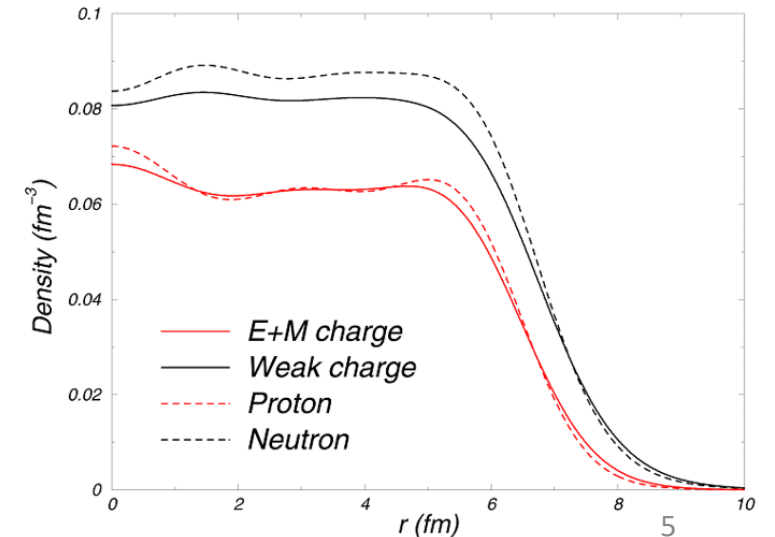


<https://arxiv.org/pdf/1904.12269.pdf>



CAP Congress

$$F_n(Q^2) = \frac{1}{4\pi} \int d^3r j_0(qr) \rho_n(r)$$



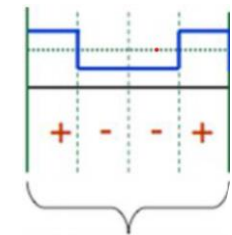
Measuring A_{PV} with ES

$$A_{PV}$$

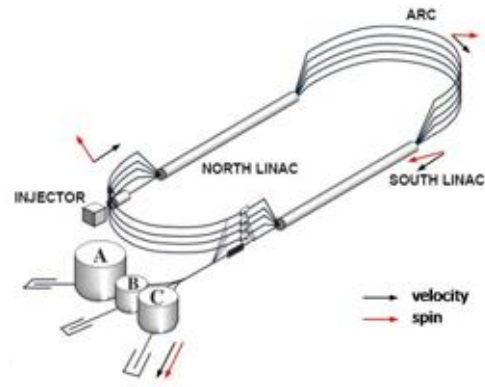
$$A_{sig} = \frac{A_{corr} - A_{back} f_{back}}{f_{sig}}$$

$$A_{corr} = A_{meas} - \sum_{i=1}^N \frac{1}{2Y} \left(\frac{\partial Y}{\partial P_i} \right) \Delta P_i$$

$$\text{where } \Delta P_i = P_+ - P_-$$



$$A_{meas} = \frac{Y_+ - Y_-}{Y_+ + Y_-}$$



Collimators

Target

spectrometer

Collimators

Magnets

detectors

main detector

tracking
detectors

background
determination

luminosity
monitors

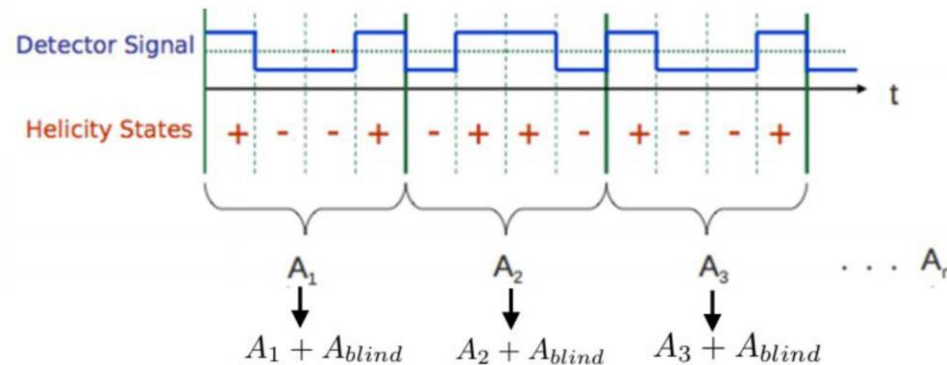
electronics

polarized
beam

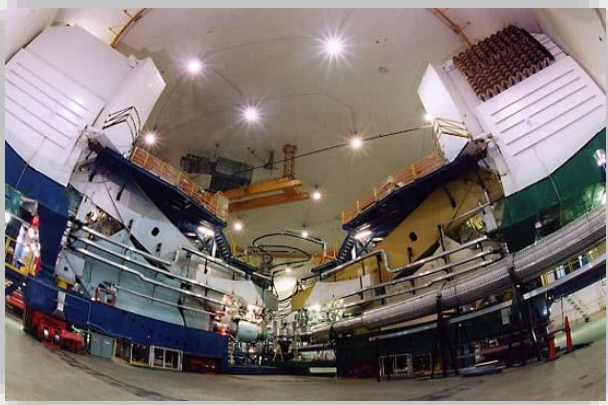
target

cooling

position



Hall A High resolution spectrometers

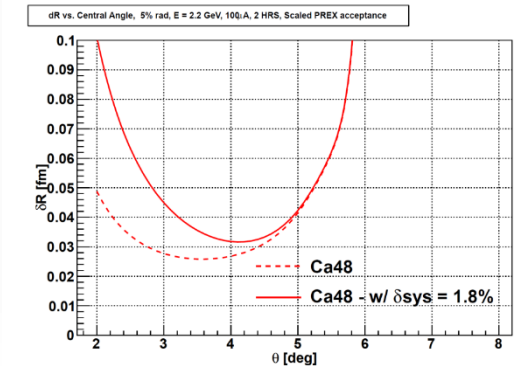
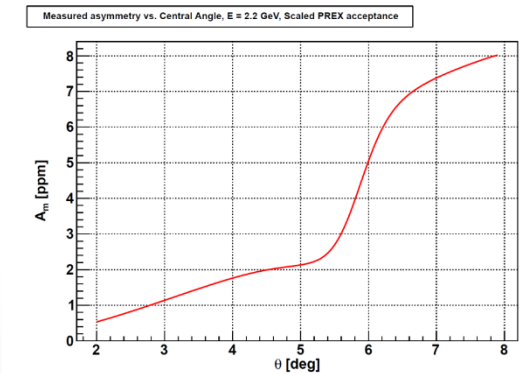
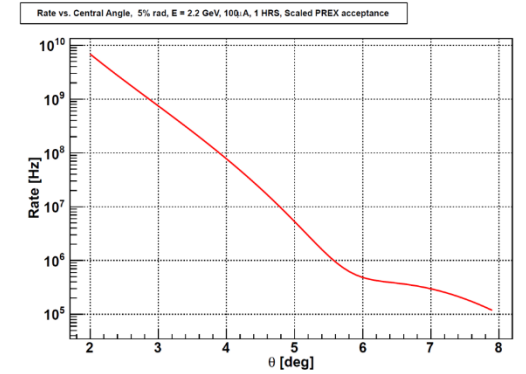
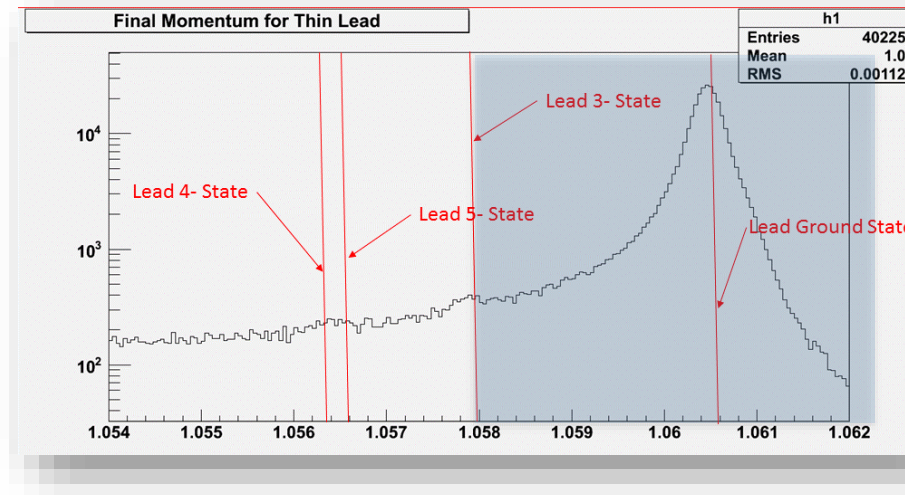
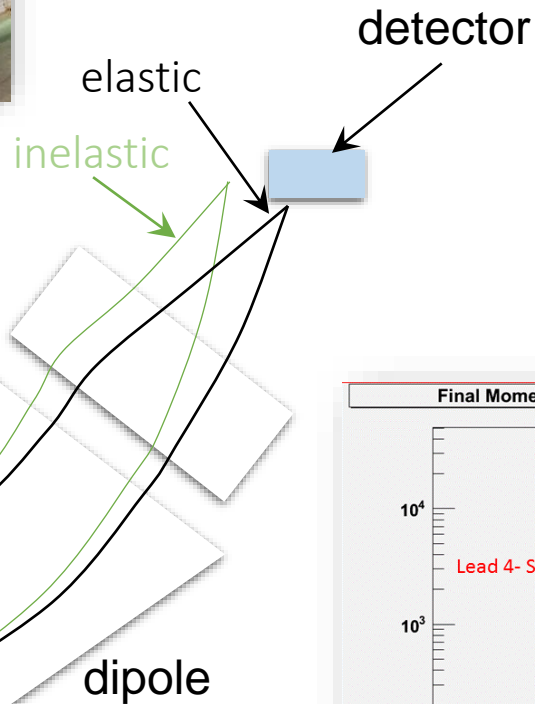


1 (2.2) GeV electron beam, 50-70 (150) μA
 high polarization, $\sim 89\%$
 helicity reversal at 120 Hz or 240 Hz

PREX (CREX)
 Parameters

0.5 (5) mm thick Pb (Ca) target
 5° (4°) scattered electrons
 $Q^2 = 0.0088$ (0.022) GeV^2/c^2
 thick and thin quartz detectors

$$A_{\text{meas}} = \frac{Y_+ - Y_-}{Y_+ + Y_-}$$



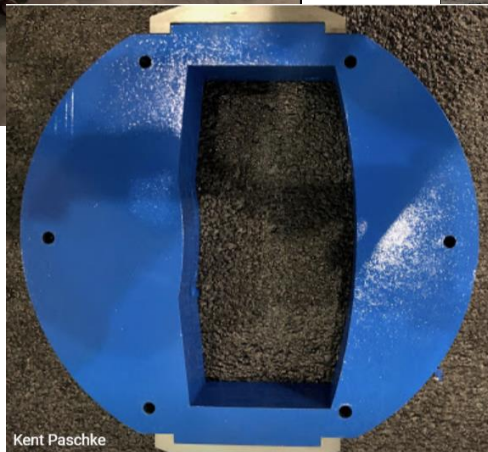
Special equipment



Septum magnet needed because to reach the low angles

Vacuum vessel to transport scattered electrons in vacuum to detector hut

Precision collimators to define the acceptance

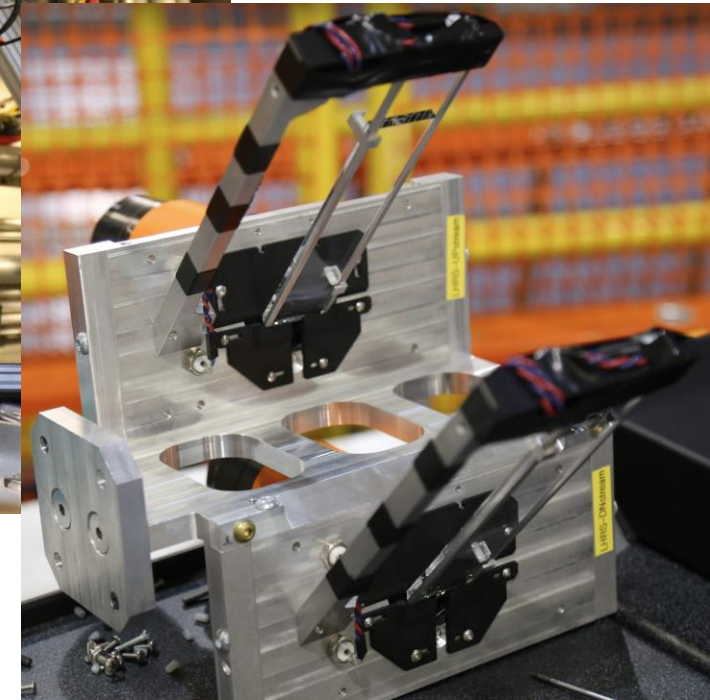


Special equipment, cont.

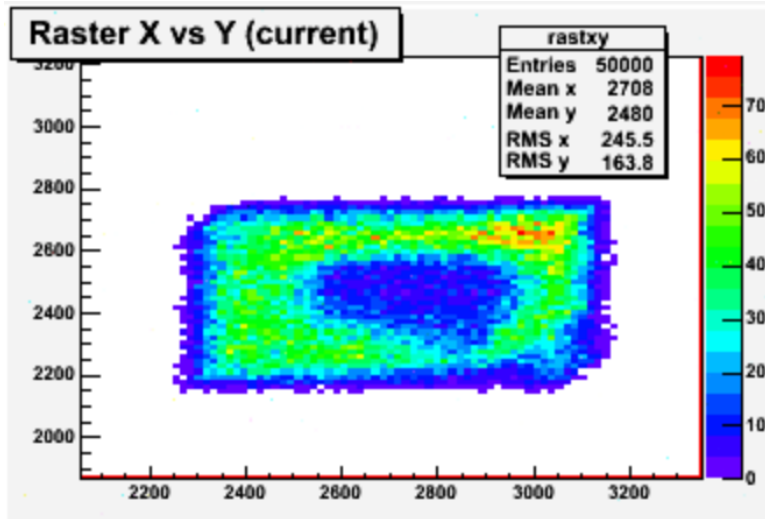


Integrating detectors (reduce
deadtime effects)

Thick and thin quartz bars
(different systematics)



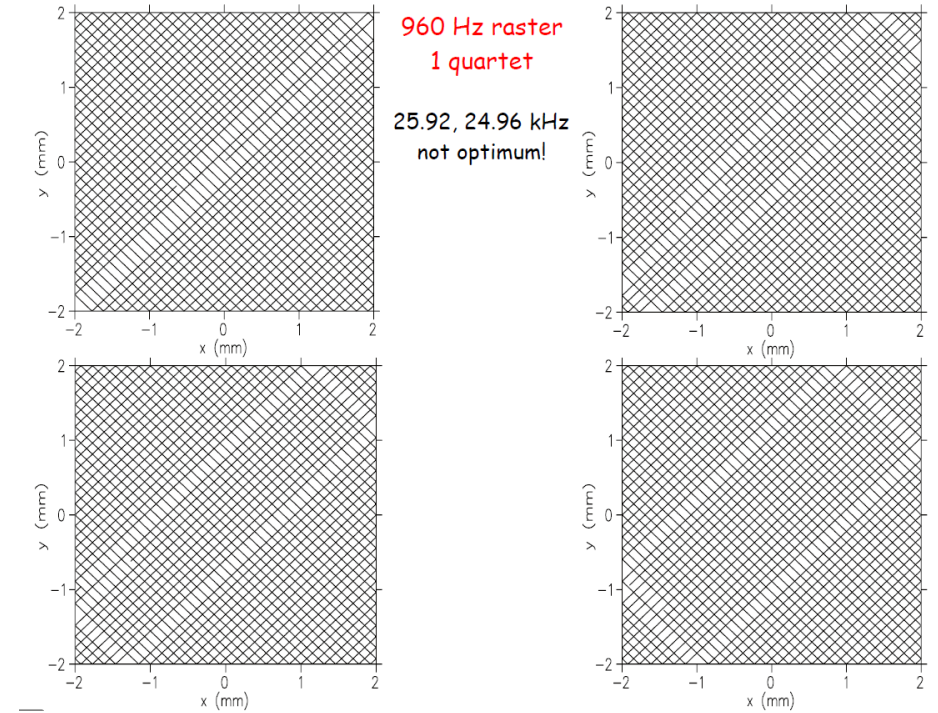
Target performance



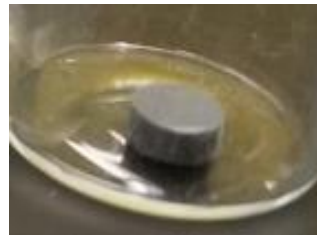
Solutions:

Sync the raster
Run with 10 targets

Acquire new ^{48}Ca
Don't expose it to air



sanded



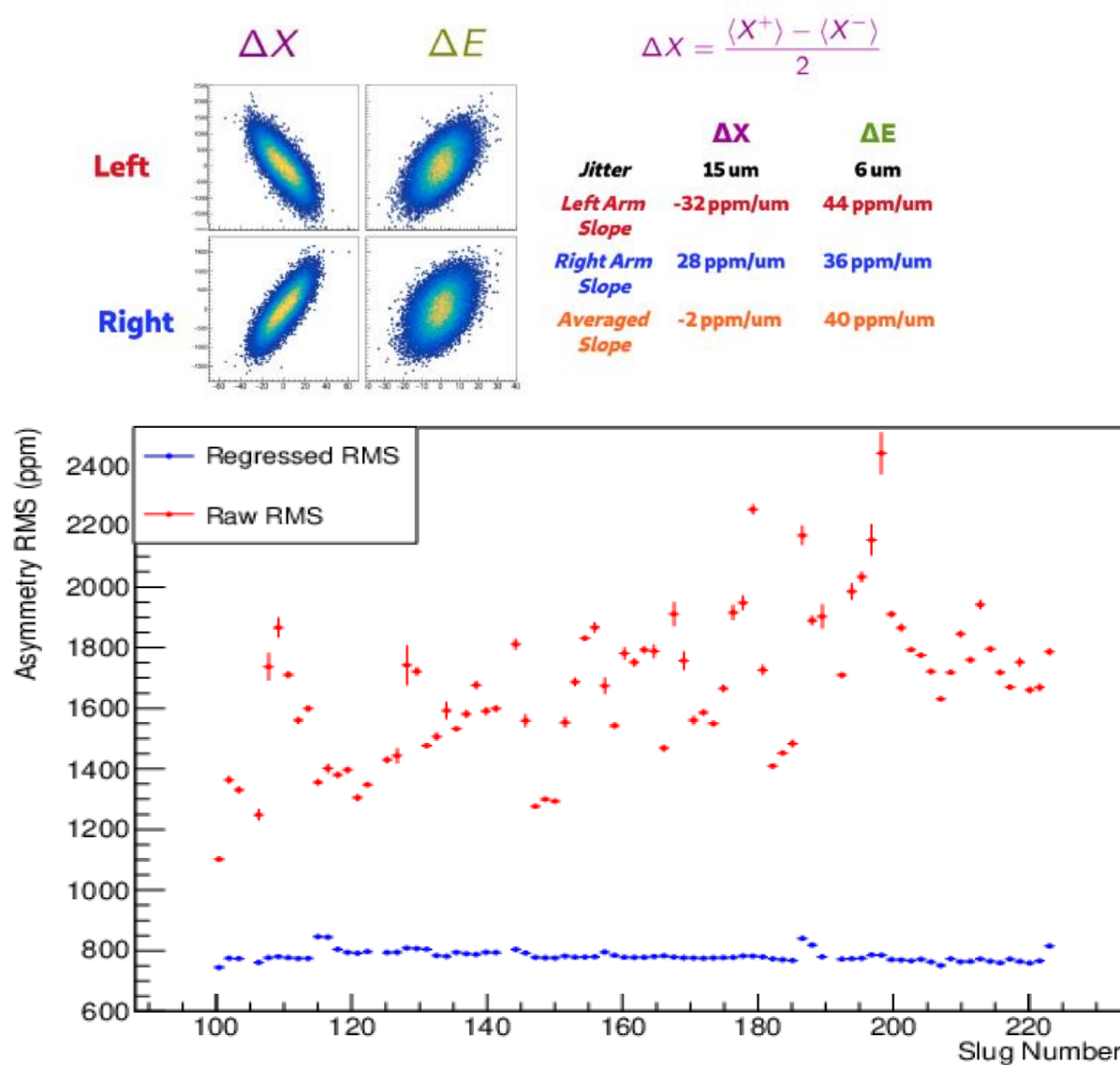
oxidized 24 hours

Calcium target for CREX was in the scattering chamber during PREX-2

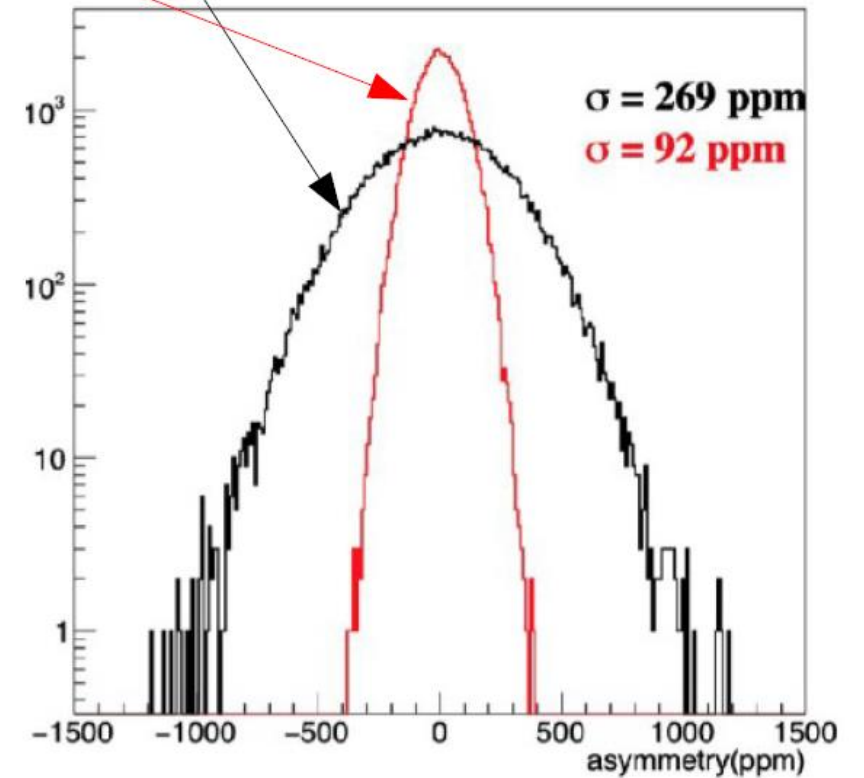
Vacuum level in target chamber monitored VERY closely



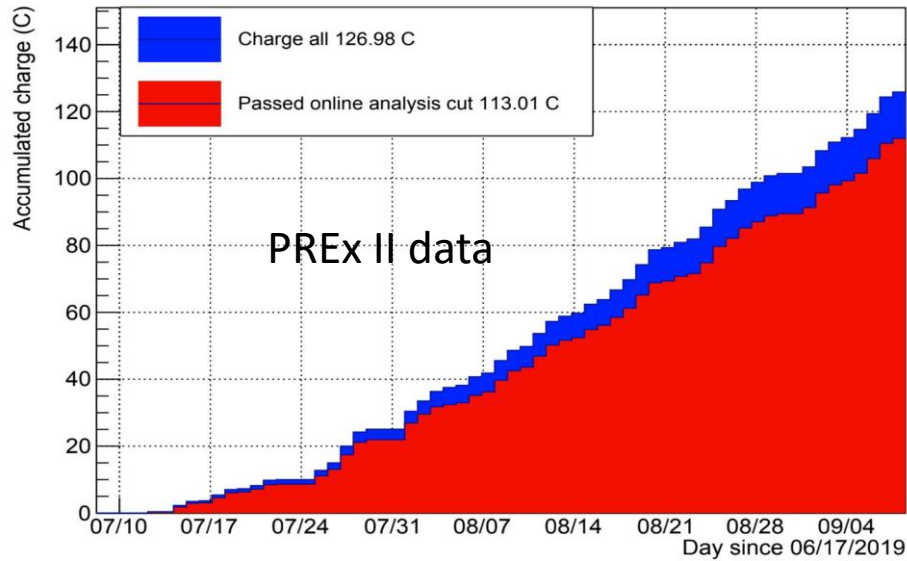
Beam Corrections, examples



$$A_{corr} = A_{meas} - \sum_{i=1}^N B_i \Delta x_i$$



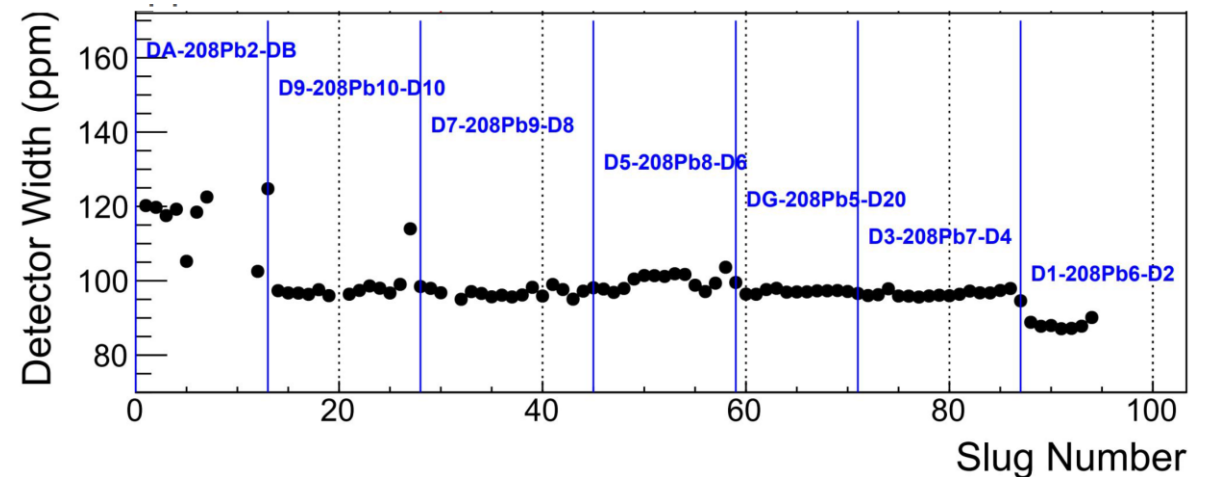
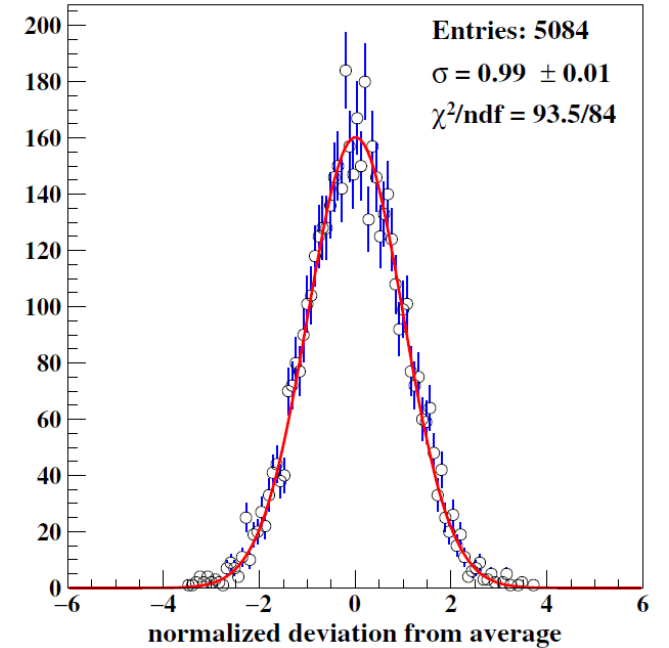
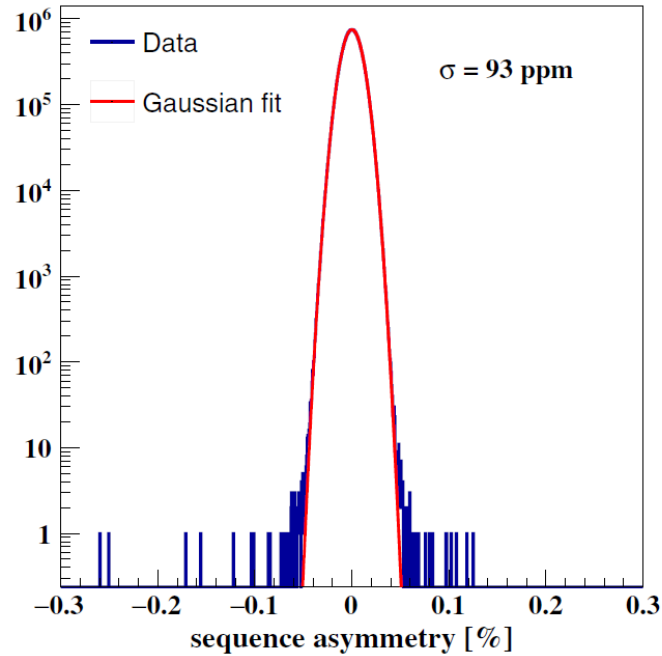
Summary of Data Quality



HWP/Wien	A_{corr}	sign	$A_{\text{PV}}^{\text{meas}}$ [ppb]	χ^2	#slugs
IN/Left	—	—	540.7 ± 29.9	46.9	27
OUT/Left	+	+	598.8 ± 29.1	31.6	29
IN/Right	+	+	506.2 ± 34.1	18.3	19
OUT/Right	—	—	536.4 ± 37.7	16.0	21

The detector system performed well; able to take ~2.5GHz on 10 x 3.5 cm² quartz in each arm

$$\sigma_A = \left(\frac{1}{\text{flip rate}} \times I (\mu A) \times R (\text{Hz}/\mu A) \times \# \text{flips} \times \# \text{dets} \right)^{-1/2}$$



Accurate Determination of the Neutron Skin Thickness of Pb208 through Parity-Violation in Electron Scattering



The High Resolution Spectrometer in Hall A at Jefferson Lab used by the PREX-2 experiment. Selected for a [Viewpoint](#) in *Physics* and an Editors' Suggestion.

From the article:

[Accurate Determination of the Neutron Skin Thickness of Pb208 through Parity-Violation in Electron Scattering](#)

D. Adhikari et al. (PREX Collaboration)

Phys. Rev. Lett. **126**, 172502 (2021)

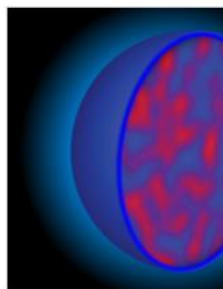
Probing the Skin

Kate Scholberg

Physics Department, Duke University, Durham

April 27, 2021 • Physics 14, 58

Researchers make the most precise measurement yet of the thickness of the neutron "skin" of a lead nucleus, with implications for the structure of neutron stars.



Primary detector element
Image: The PREX Collaboration



An experimental hall in the Thomas Jefferson National Accelerator Facility, showing a large circular structure and a person standing in the center.

APS/Alan Stonebraker

Figure 1: A cartoon image of a lead-208 nucleus, showing the mixed proton-neutron core and the neutron "skin" (left). Measuring the thickness of the neutron skin offers clues about how neutron stars are structured (right).

Physics 'What

First direct 'neutron skin' measurement

By Zack Fishman

April 27, 2021



Isaac Schmitt
4/27/21 3:16

The Academic Times



Highly Accurate Measurements Show Neutron Star "Skin" Is Less Than a Millionth of a Nanometer Thick

TOPICS: Astrophysics DOE Neutron Star Particle Physics Popular

By THOMAS JEFFERSON NATIONAL ACCELERATOR FACILITY APRIL 27, 2021



Illustration of a powerful X-ray burst erupts from a magnetar — a supermagnetized version of a stellar remnant known as a neutron star. Credit: NASA's Goddard Space Flight Center/Chris Smith (USRA)

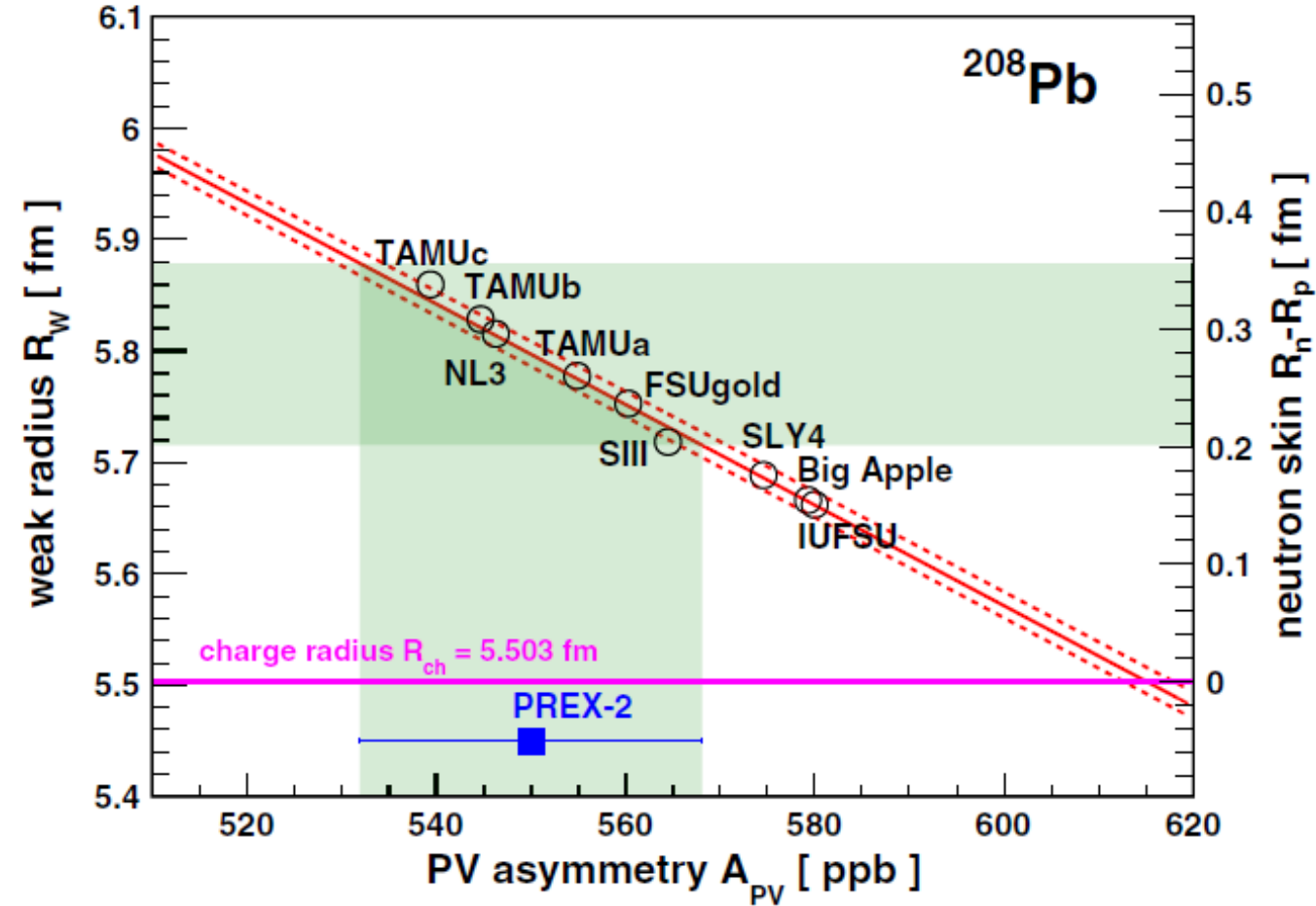
Nuclear physicists make new, high-precision measurement of the layer of neutrons that encompass the lead nucleus, revealing new information about neutron stars.

PREX-2 Result

Correction	Absolute [ppb]	Relative [%]
Beam asymmetry	-60.4 ± 3.0	11.0 ± 0.5
Charge correction	20.7 ± 0.2	3.8 ± 0.0
Beam polarization	56.8 ± 5.2	10.3 ± 1.0
Target diamond foils	0.7 ± 1.4	0.1 ± 0.3
Spectrometer rescattering	0.0 ± 0.1	0.0 ± 0.0
Inelastic contributions	0.0 ± 0.1	0.0 ± 0.0
Transverse asymmetry	0.0 ± 0.3	0.0 ± 0.1
Detector nonlinearity	0.0 ± 2.7	0.0 ± 0.5
Angle determination	0.0 ± 3.5	0.0 ± 0.6
Acceptance function	0.0 ± 2.9	0.0 ± 0.5
Total correction	17.7 ± 8.2	3.2 ± 1.5
A_{PV}^{meas} and statistical error	550 ± 16	100.0 ± 2.9

$$A_{PV}^{\text{meas}} = 550 \pm 16 \text{ (stat)} \pm 8 \text{ (syst) ppb}$$

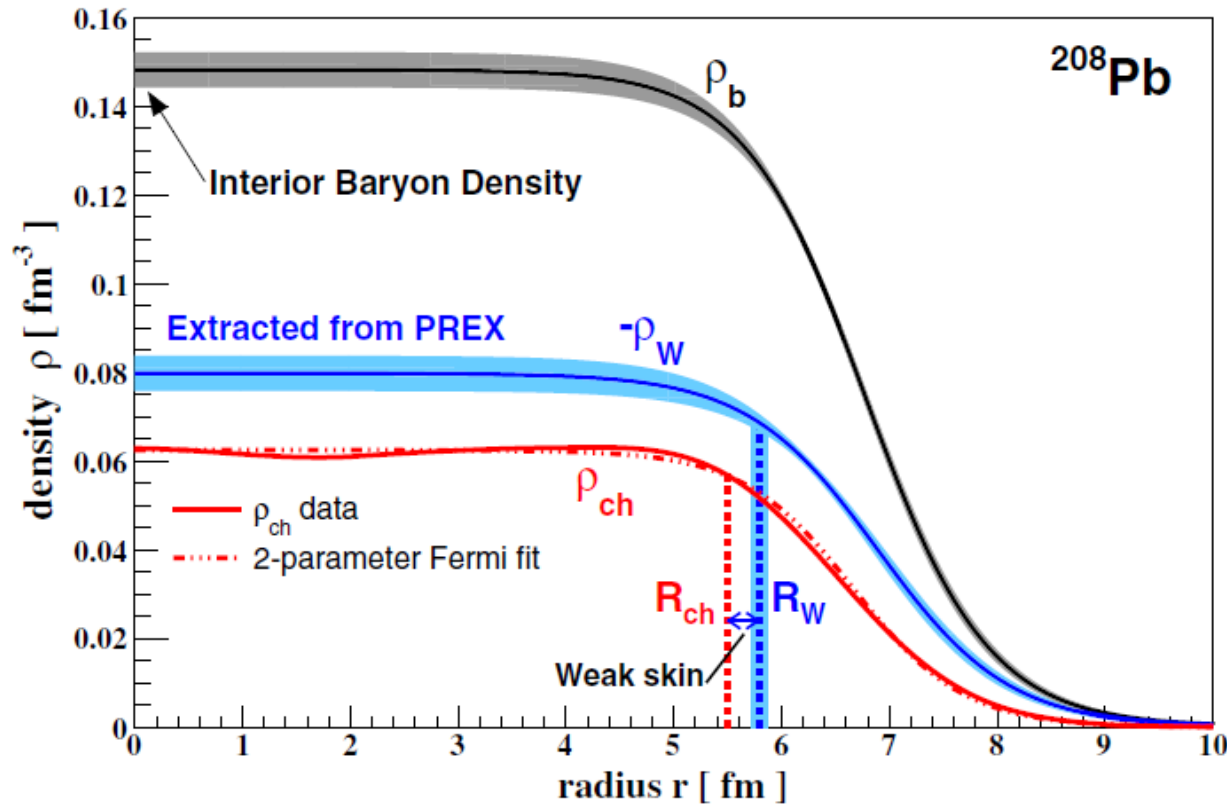
PREX 1 result: $R_n - R_p = 0.30 \pm 0.18 \text{ fm}$



$$R_w = 5.795 \pm 0.082(\text{exp}) \pm 0.013(\text{theo}) \text{ fm}$$

$$R_n - R_p = 0.278 \pm 0.078(\text{exp}) \pm 0.012(\text{theo}) \text{ fm.}$$

Interior Baryon Density



Interior weak density:

$$\rho_W^0 = -0.0798 \pm 0.0038 \text{ (exp)} \pm 0.0013 \text{ (theo)} \text{ fm}^{-3}.$$

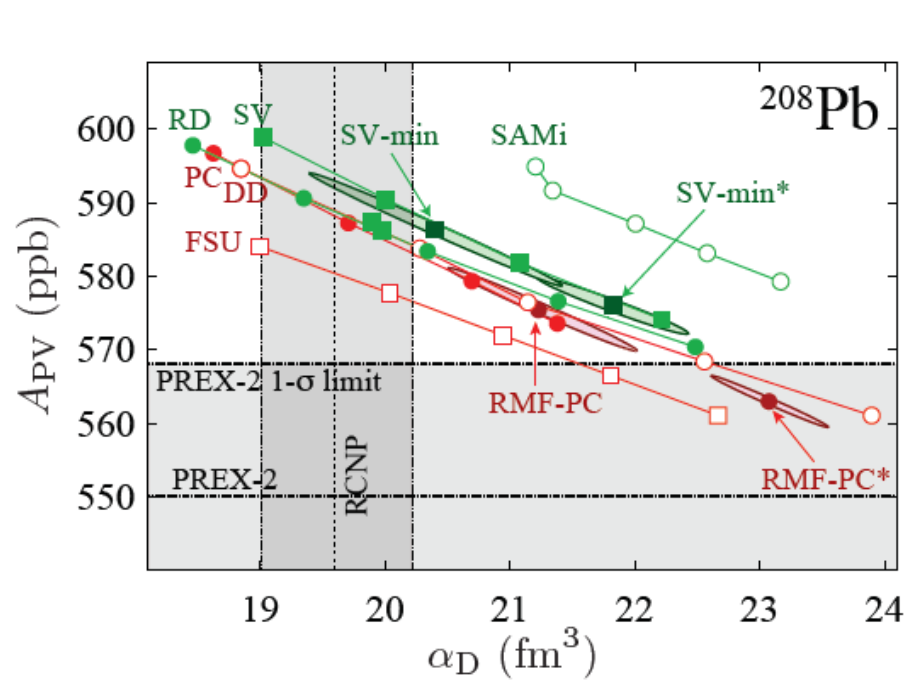
Interior baryon density from
PREX combined and charge
density data:

$$\rho_b^0 = 0.1482 \pm 0.0040 \text{ fm}^{-3}$$

Note: acceptance function
needed to interpret correctly

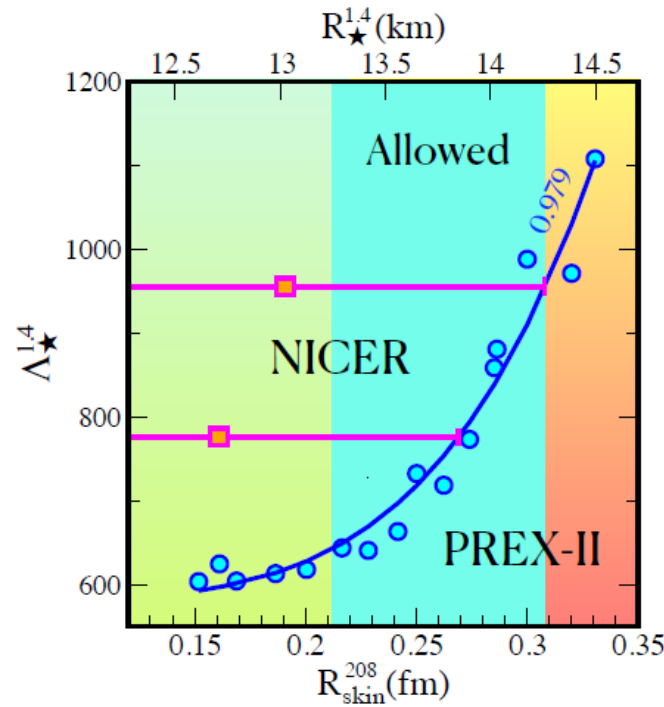
$$\langle A_{PV} \rangle = \frac{\int d\theta \sin \theta A(\theta) \frac{d\sigma}{d\Omega} \epsilon(\theta)}{\int d\theta \sin \theta \frac{d\sigma}{d\Omega} \epsilon(\theta)}$$

Context and Implications - Tension



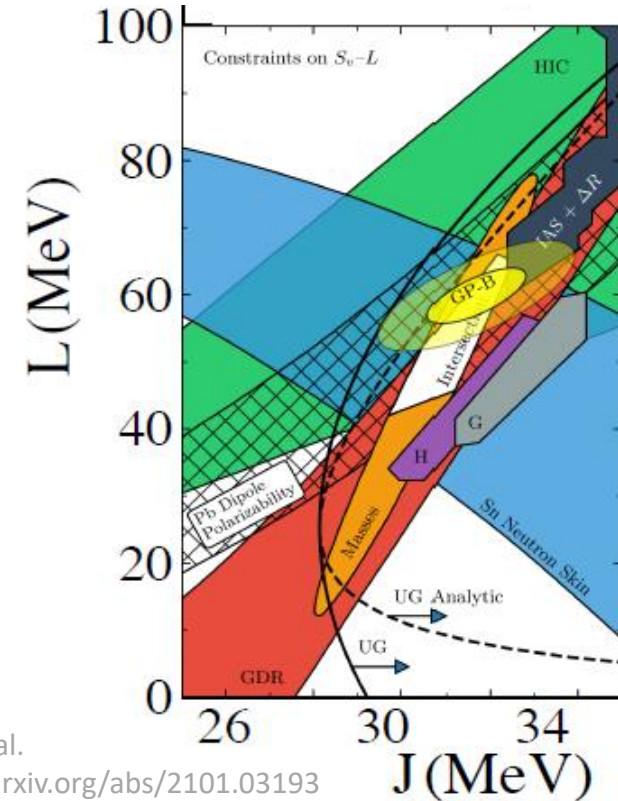
Reinhard et al.
<https://arxiv.org/pdf/2105.15050.pdf>

There is no model which is able to simultaneously reproduce A_{PV} and α_D within the experimental 1 σ error bands



$R_{\text{skin}}^{208} \lesssim 0.31$ fm and a lower limit on the stellar radius of $R_{\star}^{1.4} \gtrsim 13.25$ km

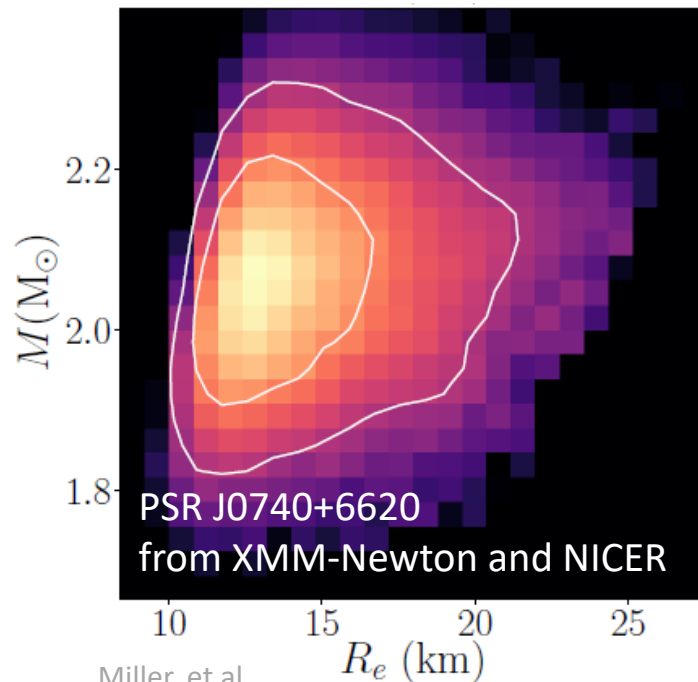
Reed et al.
<https://arxiv.org/abs/2101.03193>



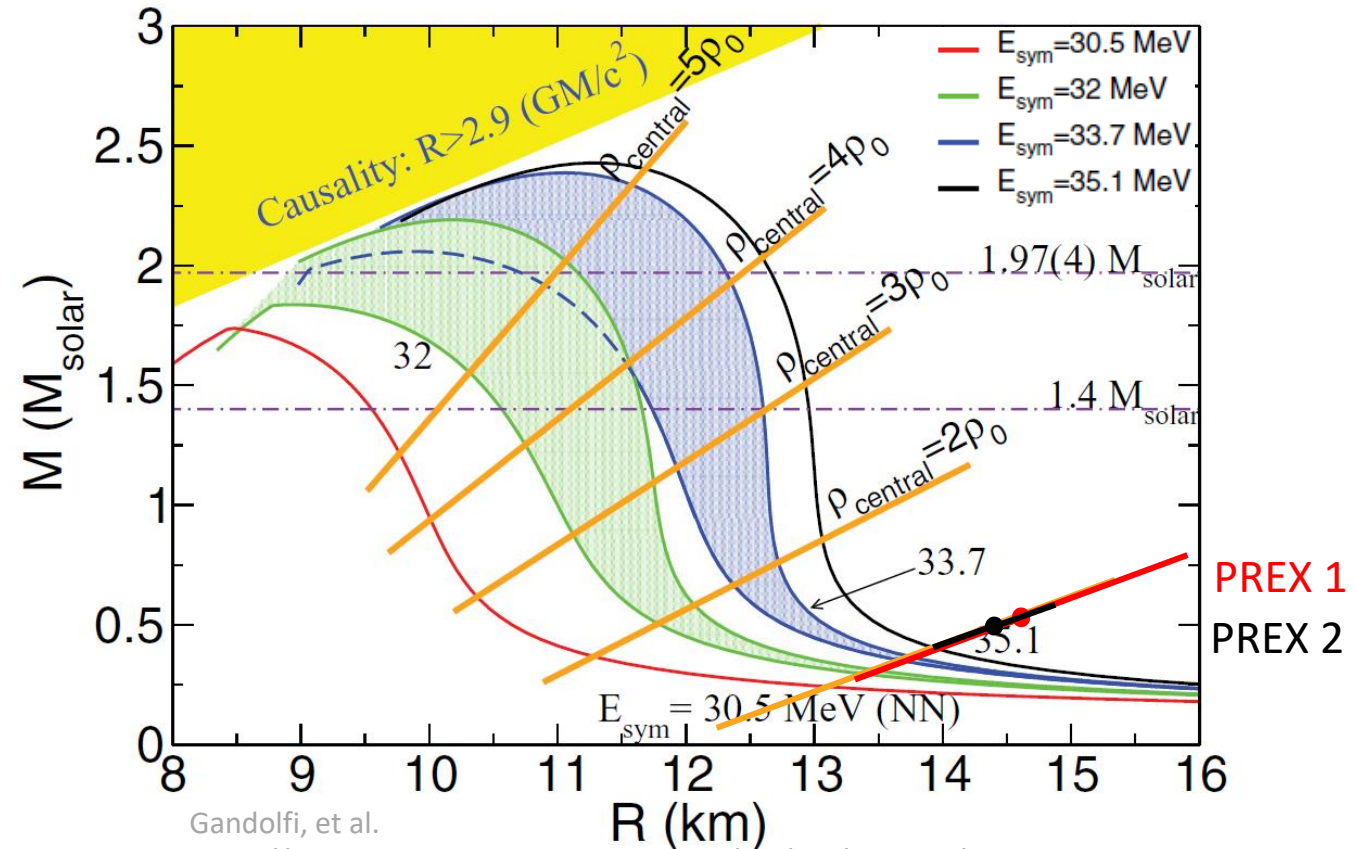
The PREX-II result is relatively model-independent

Relating PREX and CREX to Neutron Stars

- The mass of a neutron star as a function of radius can be determined from neutron star observations
- Nuclear structure models can produce EOS curves relating the mass and radius, using different assumptions (values of E_{sym} , NNN interactions)



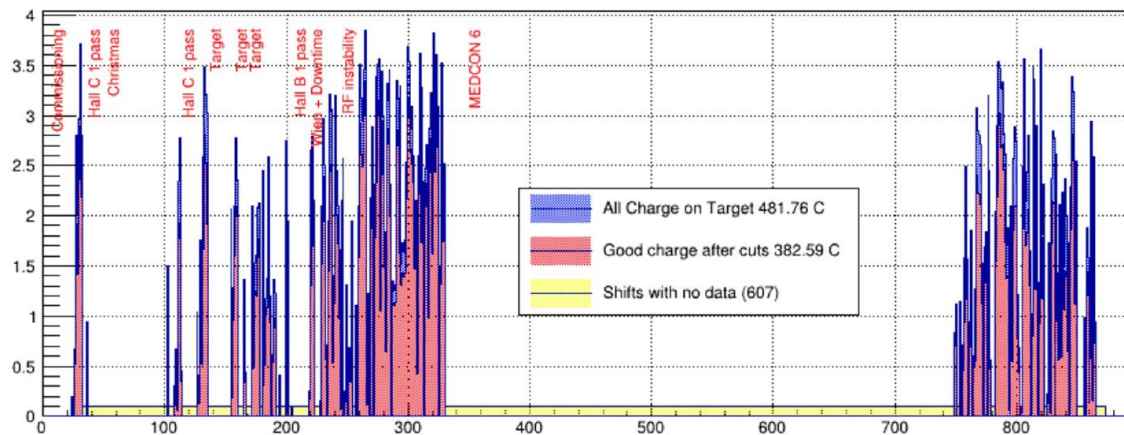
Miller, et al.
<https://arxiv.org/abs/2105.06979>



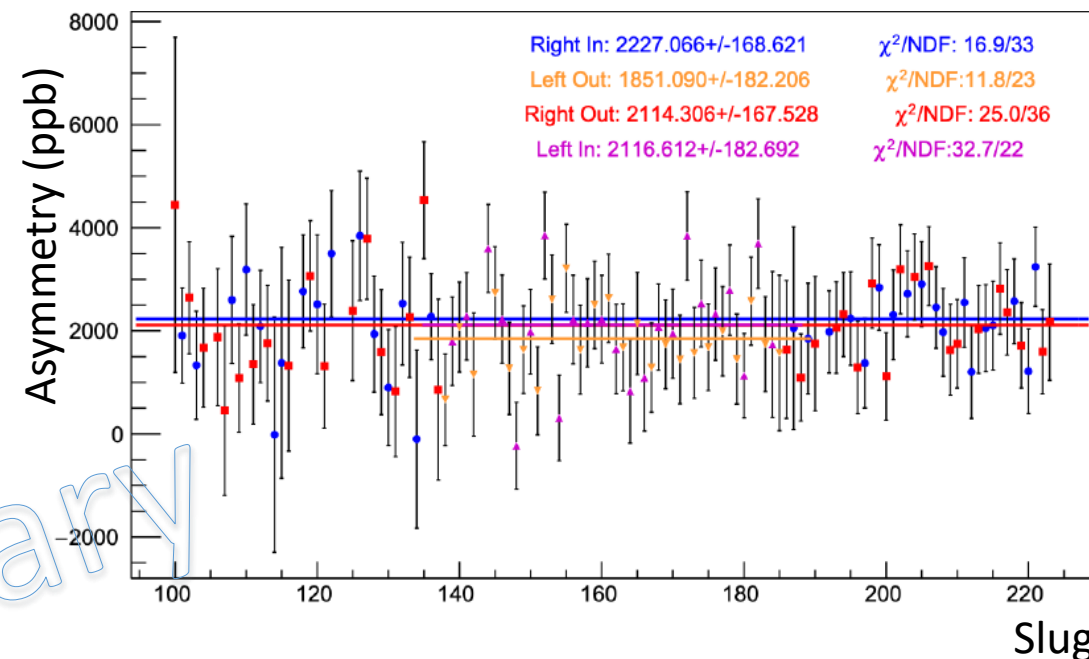
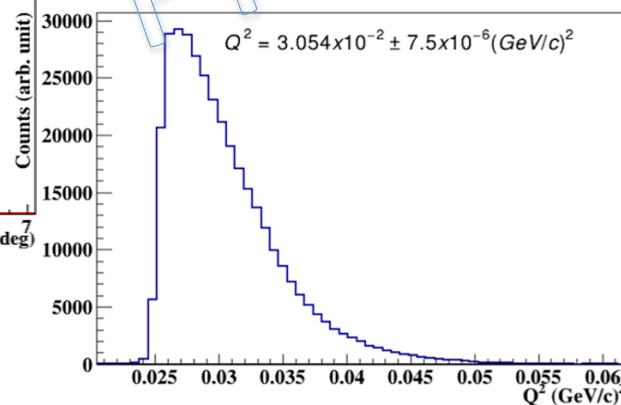
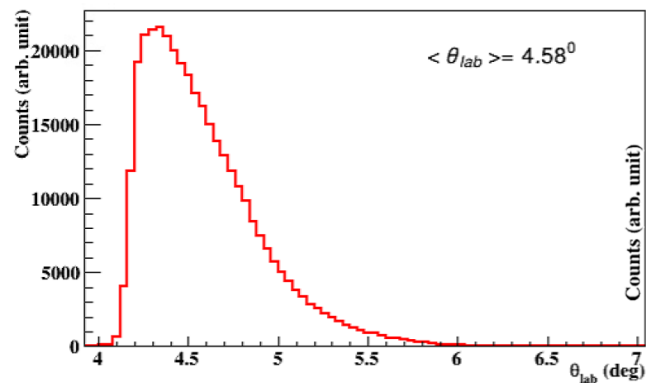
- In this plot, $E_{\text{sym}} = 35.1$ MeV corresponds to $L = 63.6$ MeV ($E_{\text{sym}} = 30.5$ MeV $\sim L = 31.1$ MeV)
- Combined PREX gives 106 ± 37 MeV (in a different model)
- CREX will help narrow the width of the 3N bands

CREX Status

Charge total vs shift



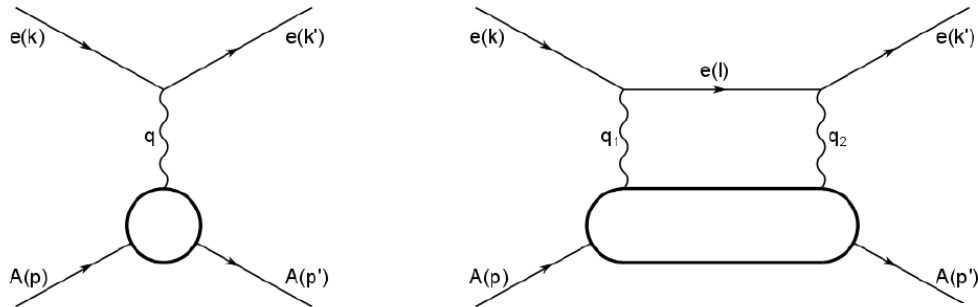
Data collected at beam energy of 2.2 GeV @150 μ A collecting a total of 481 C during production running



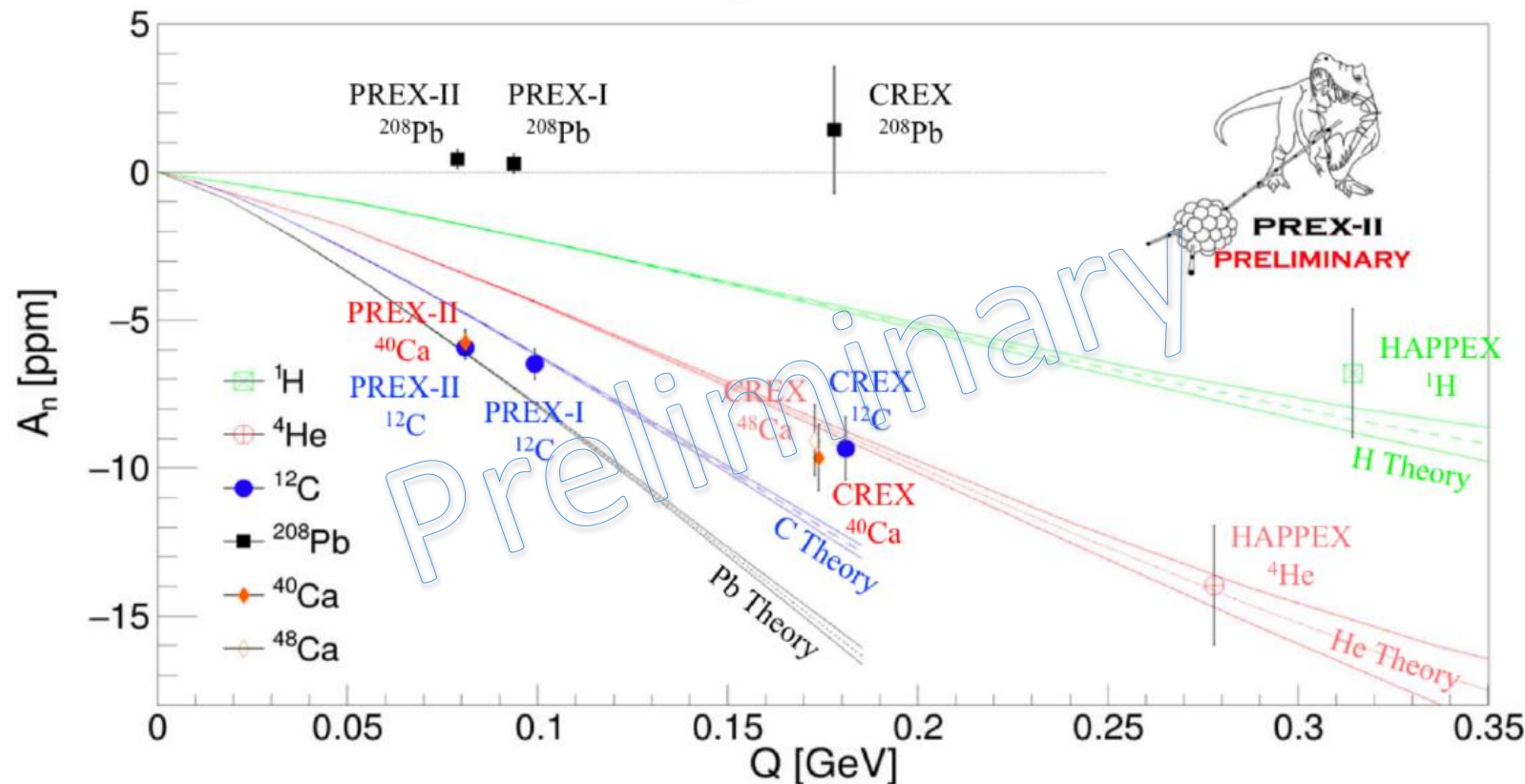
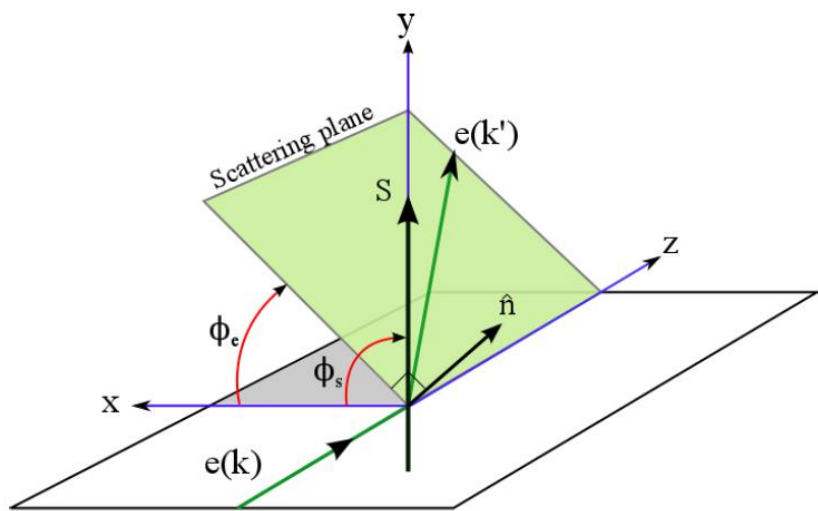
- CREX data is being analyzed
- Expect to unblind by end of summer
- Publish this fall

Expected systematic	
Charge Normalization	0.1%
Beam Asymmetries	0.3%
Detector Non-linearity	0.3%
Transverse	0.1%
Polarization	0.8%
Inelastic Contribution	0.2%
Effective Q^2	0.8%
Total	1.2%

Transverse

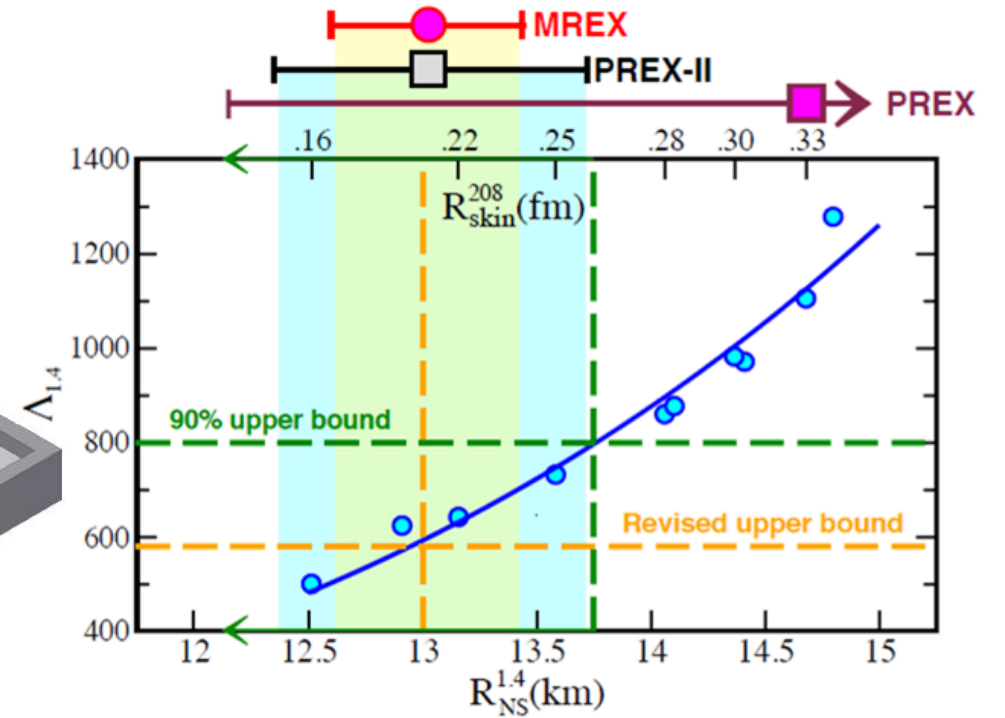
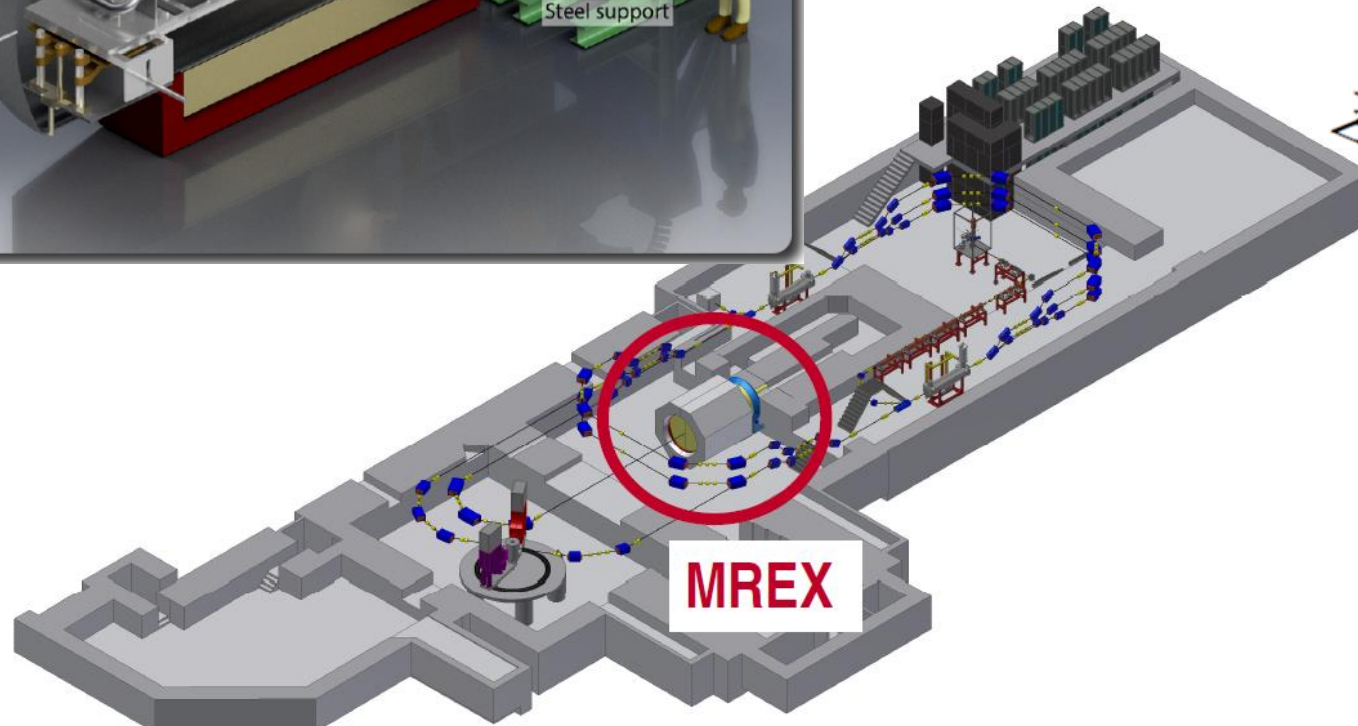
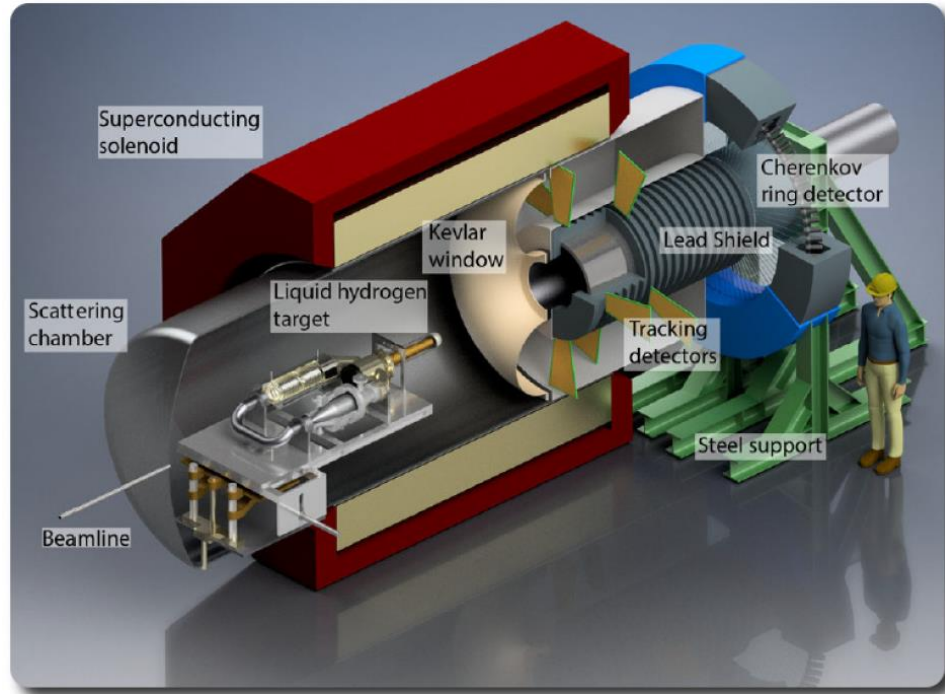


$$A_n = \frac{\sigma_{\uparrow} - \sigma_{\downarrow}}{\sigma_{\uparrow} + \sigma_{\downarrow}} \approx 0 + \frac{2\text{Im}(T_{2\gamma} \cdot T_{1\gamma}^*)}{|T_{1\gamma}|^2}$$



$$\mathcal{A}_{meas}(\phi_e) = \mathcal{A}_n(\vec{S}_e \cdot \hat{n}) = \mathcal{A}_n |\vec{S}_e| \sin(\phi_e - \phi_s)$$

MREX – at the Mainz Microtron - MESA



<https://arxiv.org/pdf/1904.12269.pdf>

Conclusion

- PREX 2 is complete

- Weak form factor
- Weak radius
- Neutron skin
- Interior baryon density

$$L = 106 \pm 37 \text{ MeV}$$

In tension with NICER,
dipole polarizability, tidal
deformability

- CREX data is being analyzed – expect to publish in Fall 2021
- Future experiment at Mainz to improve the precision